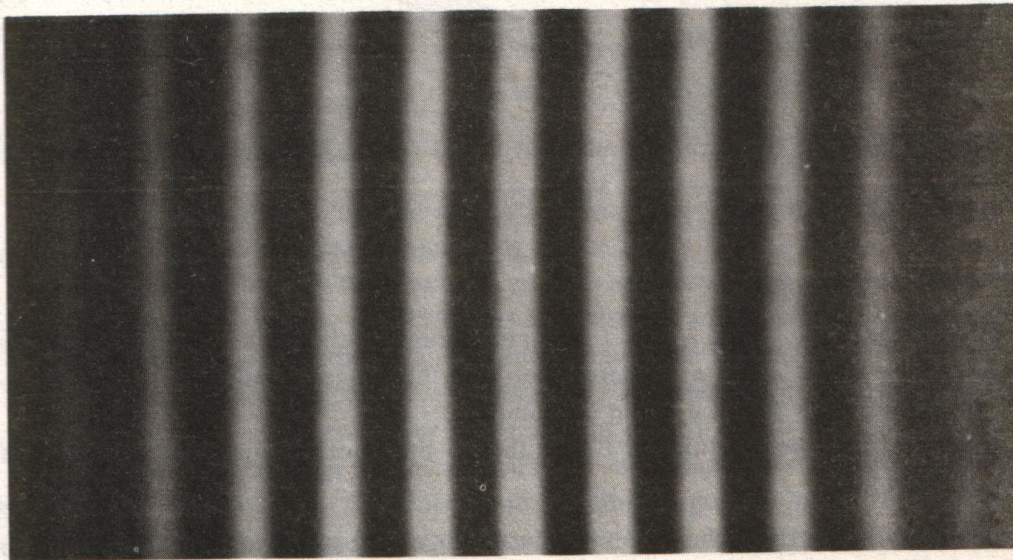
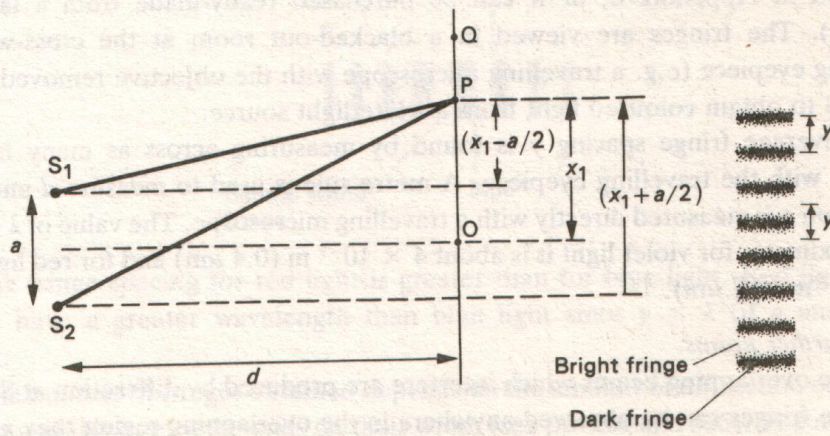


(a)



(b)

Fig. 8.3



(a)

(b)

Fig. 8.4

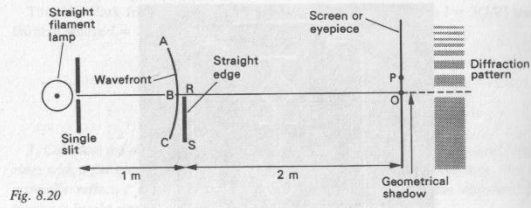


Fig. 8.20

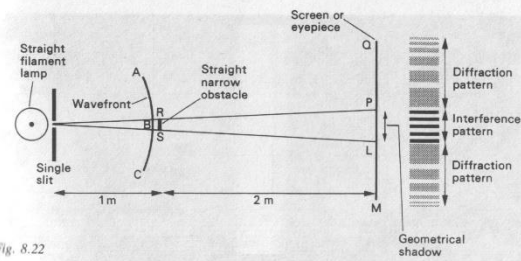
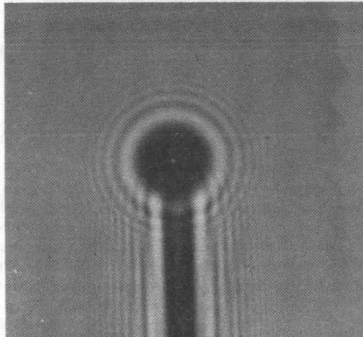


Fig. 8.22

length parallel to the filament. Initially the slit is opened wide and the lens moved to give a sharp image of the filament on the translucent screen. The shield and the large stop cut out stray light.

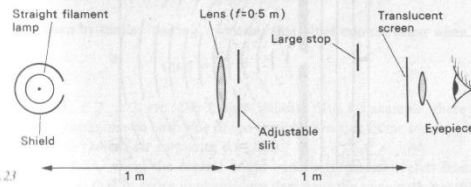


Fig. 8.23

The diffraction pattern depends on the slit width and is observed from behind the

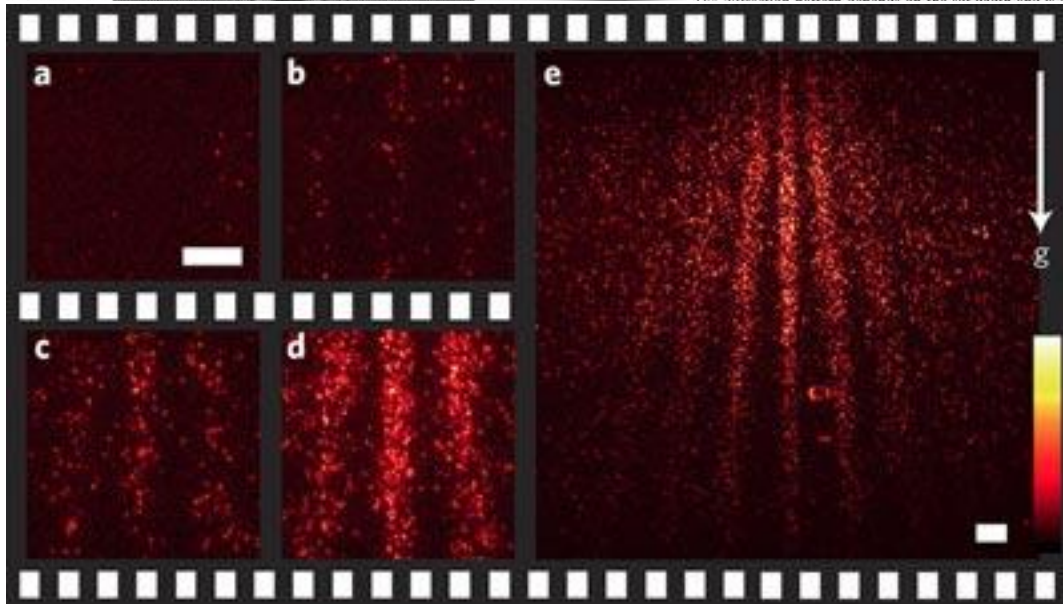


Table 22.1 The electromagnetic spectrum

Type	Frequency	Wavelength	How made	Uses	Photon energy
long-wave radio	~250 kHz	~1200 m	oscillating currents in aerials	radio	~10 ⁻³⁸ J
medium-wave radio	~1000 kHz (1 MHz)	~300 m	oscillating currents in aerials	radio	~10 ⁻²⁷ J
short-wave radio	~10 MHz	~30 m	oscillating currents in aerials	radio	~10 ⁻²⁵ J
VHF	~100 MHz	~3 m	oscillating currents in aerials	radio	~10 ⁻²⁵ J
UHF	~400 MHz	~1 m	oscillating currents in aerials	television	~10 ⁻²⁵ J
microwaves	~2.5 GHz	~10 cm	directly produced in waveguides	radar, cooking, communicating	~10 ⁻²⁴ J
infra-red	~10 ¹⁴ Hz	~1 μm (> 700 nm)	hot bodies, LEDs	night-sights, heating, short-distance communication	~10 ⁻¹⁹ J ~1 eV
visible	~5 × 10 ¹⁴ Hz	700 – 400 nm	very hot bodies, LEDs	seeing, etc	2 – 3 eV
ultra-violet	>7.5 × 10 ¹⁴ Hz	<400 nm	extremely hot bodies, sparks, discharge tubes	sun-tanning, detecting invisible marking, sterilising	> 3 eV
X-rays	~10 ¹⁸ Hz	~10 ⁻¹⁰ m	stopping fast electrons	X-raying people and materials	~10 000 eV
gamma rays (overlap with X-rays)	~10 ²⁰ Hz	~10 ⁻¹² m	nuclear decay	X-raying thick objects, killing cancerous cells, sterilising	~1 MeV ~10 ⁻¹³ J
cosmic rays	very high	very short	from distant parts of the Universe	just cause a hazard	up to tens of joules



radiowaves



radio transmitter




radio



television

microwaves



microwave oven



microwave transmitter

infrared



electric fire



infra-red grill

ROYGBIV

visible




sight




photography

ultraviolet




sun bed




security

X-rays



Gamma-rays



Section D Waves

D1(T) Investigating oscillations

When a mass hanging from a vertical spring is pulled down and released it bobs up and down. Studying this simple arrangement gives a clear picture of the essential features of any oscillating system.

1 Does the time period depend on the amplitude of the oscillations?

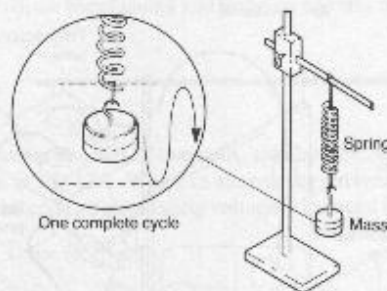


Fig 1A Timing oscillations

Displace a mass downwards from equilibrium by a measured distance and release it (see Fig 1A). Time 20 complete oscillations and work out the time period, T . Repeat the measurement several times. Estimate and record the probable errors in your measurements.

Repeat this for different initial displacements. Record all your measurements in a table.

Question

- 1 Present your results as a graph, showing the experimental errors in each measurement on your graph. What conclusions can you draw from your results?

2 How does the displacement vary with time after the mass is released?

Use a tickertimer and tickertape to record the motion of the mass for the first half-cycle after release. Fig 1B shows one way to do this.

Use your tape to make a graph of displacement against time. From the graph, work out the time period.

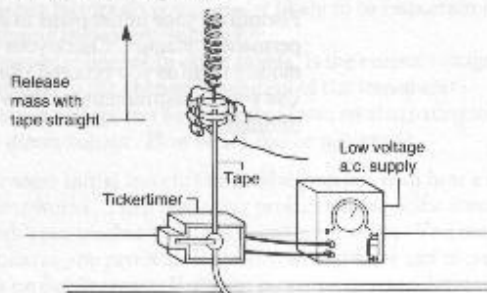


Fig 1B Recording the motion

Questions

- Sketch graphs on the same axes to show how the velocity and the acceleration vary with time. What can you say about the acceleration in relation to the displacement?
- Describe the motion of the mass after release. The amplitude gradually becomes smaller and smaller. Why?

How does the time period depend on the mass?

Measure the time period for different masses hanging from the spring. Record your measurements and estimate the probable error in them.

Questions

- The time period, T may be worked out from the equation:

$$T = 2\pi\sqrt{\frac{m}{k}}$$

where m = mass hung on the spring and k = the spring constant.

Plot a suitable graph to confirm the above relationship and work out a value for k .

- Determine the spring constant, k , by loading the spring with different known weights and measuring the extension. Plot a graph of your measurements and work out a value for k from this graph. Compare this value of k with the value obtained above.

Q2(I) The simple pendulum

Measuring the time period, T

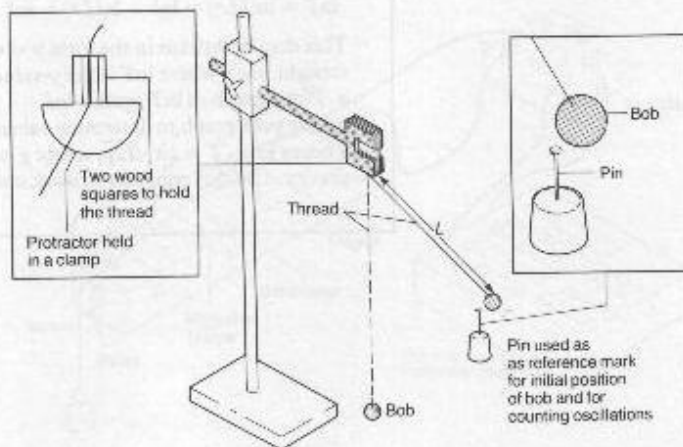


Fig 2A The simple pendulum

Set up the pendulum as in Fig 2A. Displace the pendulum by a measured angle from its equilibrium position and time 20 complete oscillations.

How reliable is your timing? Repeat the measurement ten times, displacing the pendulum by the same angle each time. Record your measurements.

Questions

- 1 Work out the average timing for 20 oscillations, T_{20} .
- 2 From the spread of measurements, estimate the probable error for T_{20} .
- 3 How many times is it necessary to make the measurements to obtain a reliable result?

2 Does the time period depend on the amplitude of oscillations?

Displace the pendulum by 5° from equilibrium and time 20 oscillations after it has been released. Repeat the timings to obtain a reliable measurement. Record your measurements and work out the time period, T .

Repeat and record the measurements for initial displacement, θ , = 10° , 15° , 20° , 25° and 30° .

Question

- 4 Plot a graph of T (on the vertical axis) against θ . Use your graph to decide if T depends on θ .

3 How does the time period depend on length?

Measure the time period, T , for at least five different lengths, L . For each timing the initial displacement should not exceed 10° . Record all your measurements in a table.

Questions

- 5 Assume the relationship between T and L is of the form $T = kL^n$, where k and n are constants. Using the theory of logarithms:

$$\ln T = \ln(kL^n) = \ln k + \ln(L^n) = \ln k + n \ln L$$

This may be written in the form $y = mx + c$ (the equation for a straight line), where $\ln T$ is the y -value and $\ln L$ is the x -value.

- a Plot a graph of $\ln T$ against $\ln L$.
 - b Use your graph to determine values for n and k .
- 6 Theory gives $T = 2\pi \sqrt{L/g}$, where g is the acceleration due to gravity. Use this equation to work out a value for g .

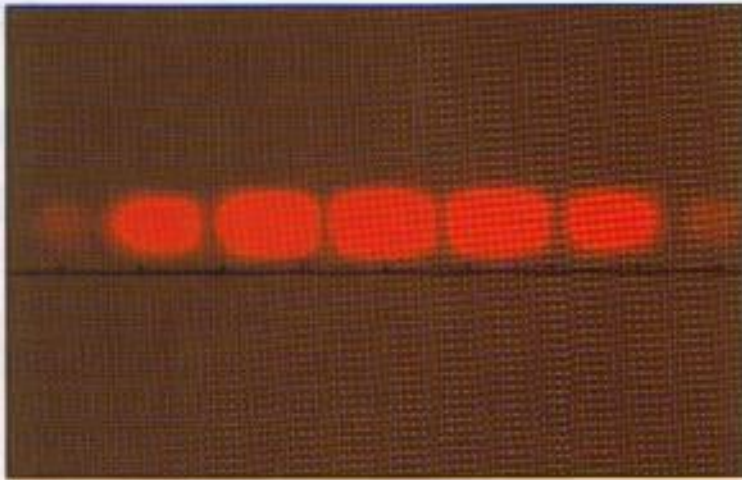


Figure 16.4 Superposition pattern for laser light through a double slit

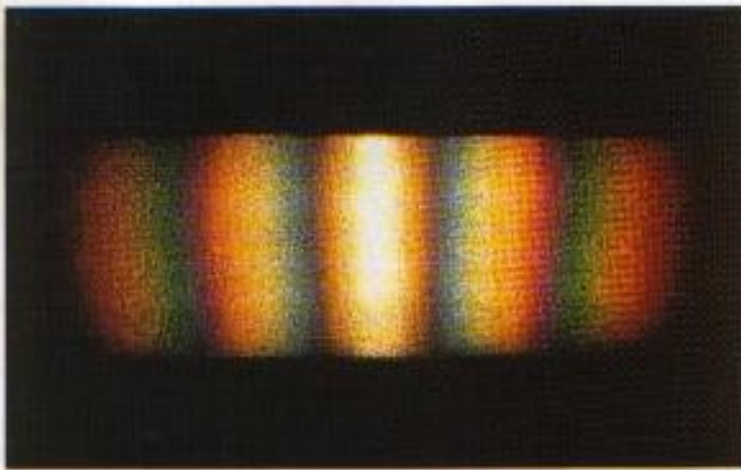


Figure 16.5 Superposition pattern caused by white light through a double slit

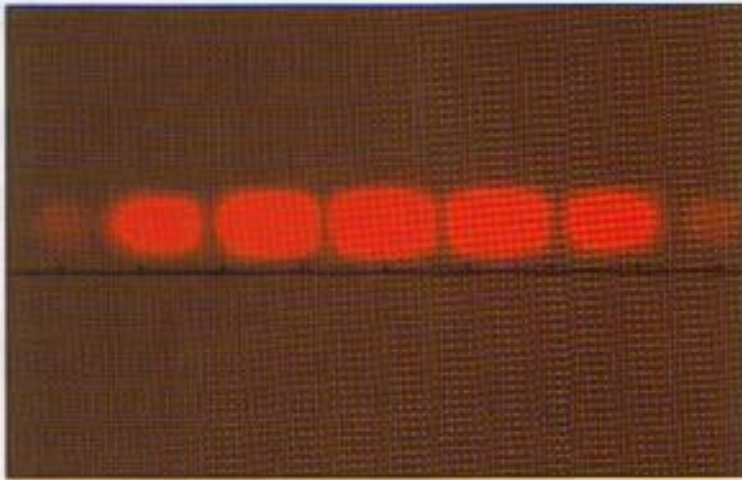


Figure 16.4 Superposition pattern for laser light through a double slit

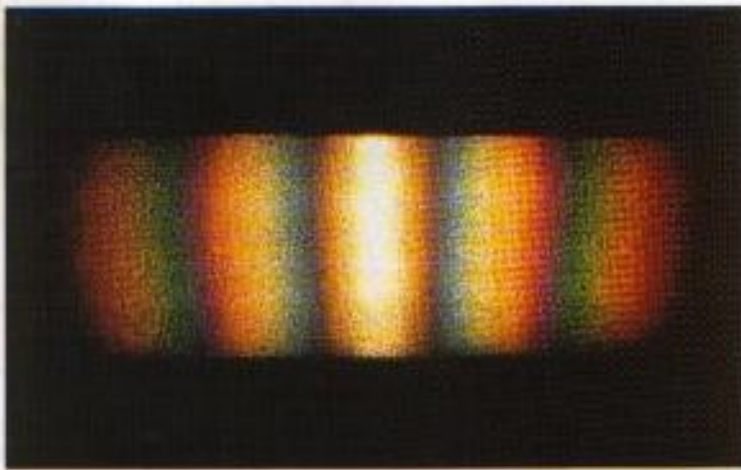
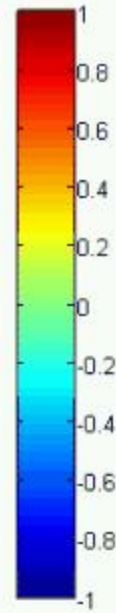
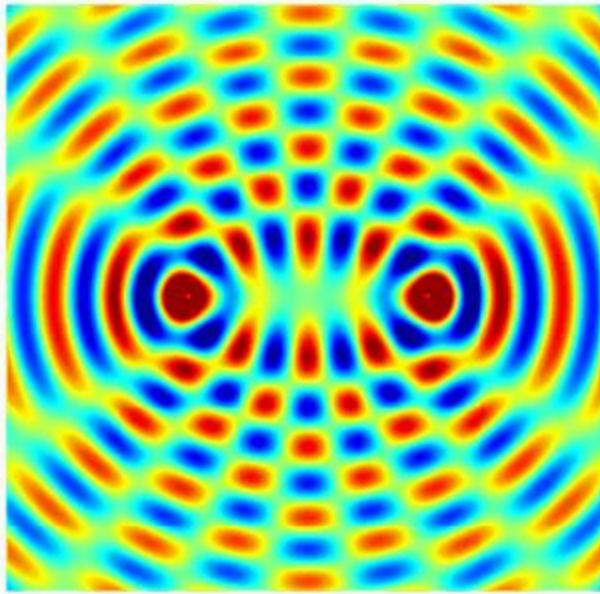


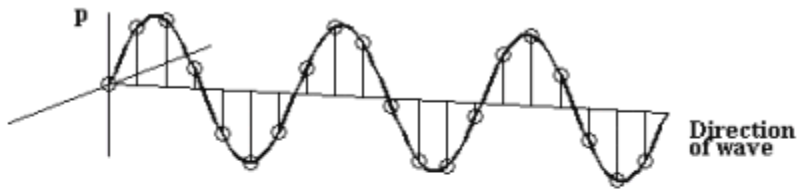
Figure 16.5 Superposition pattern caused by white light through a double slit

Four wavelengths spacing

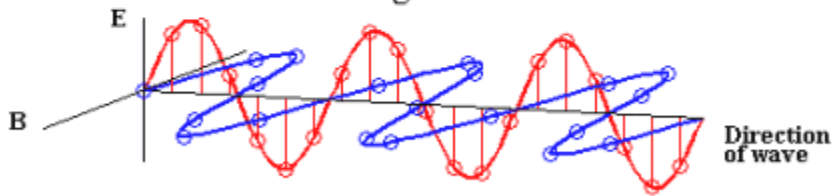


isvr

Mechanical wave

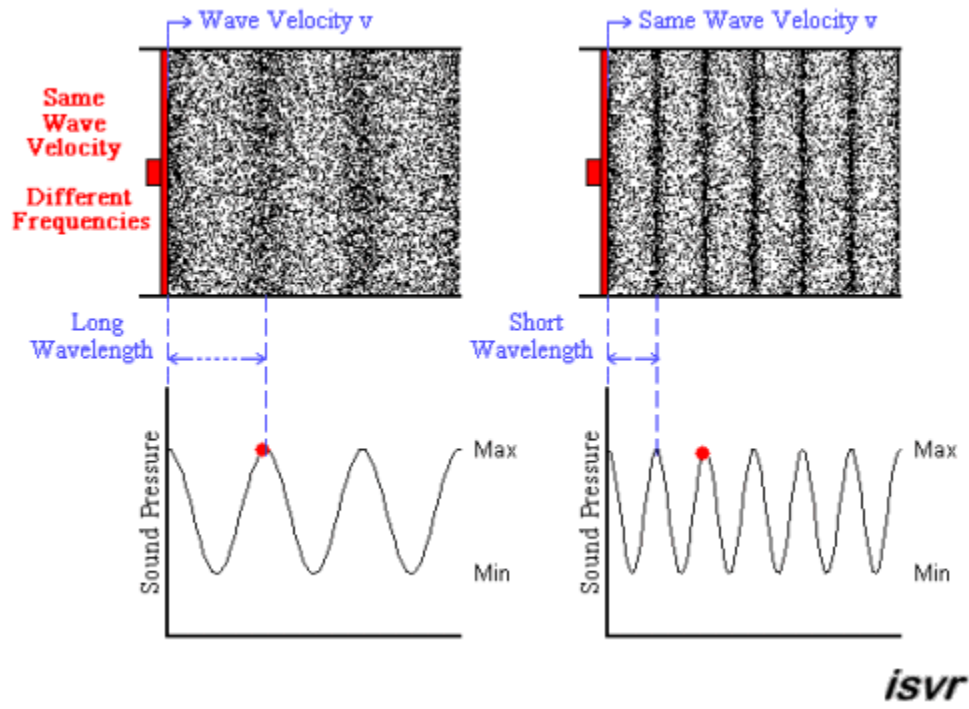
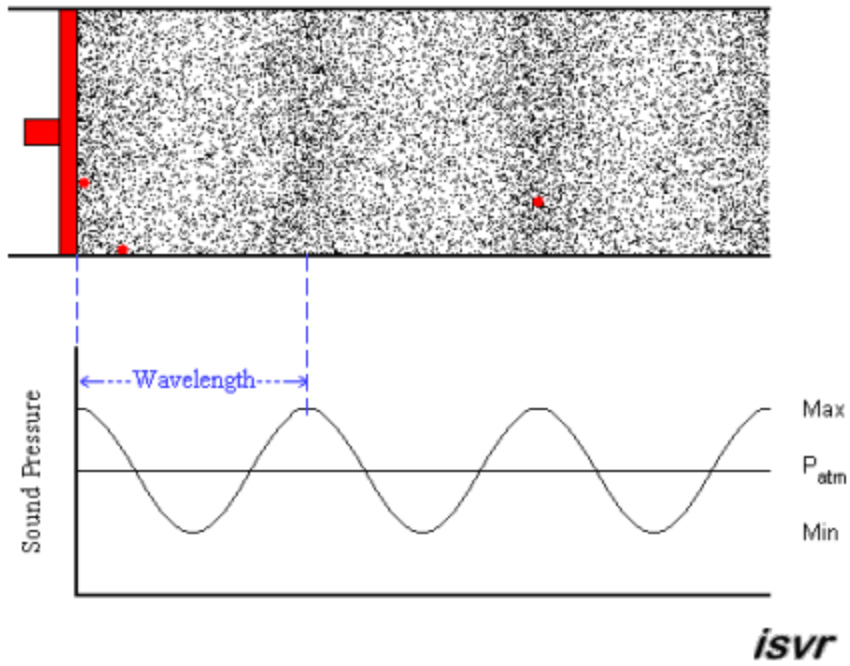


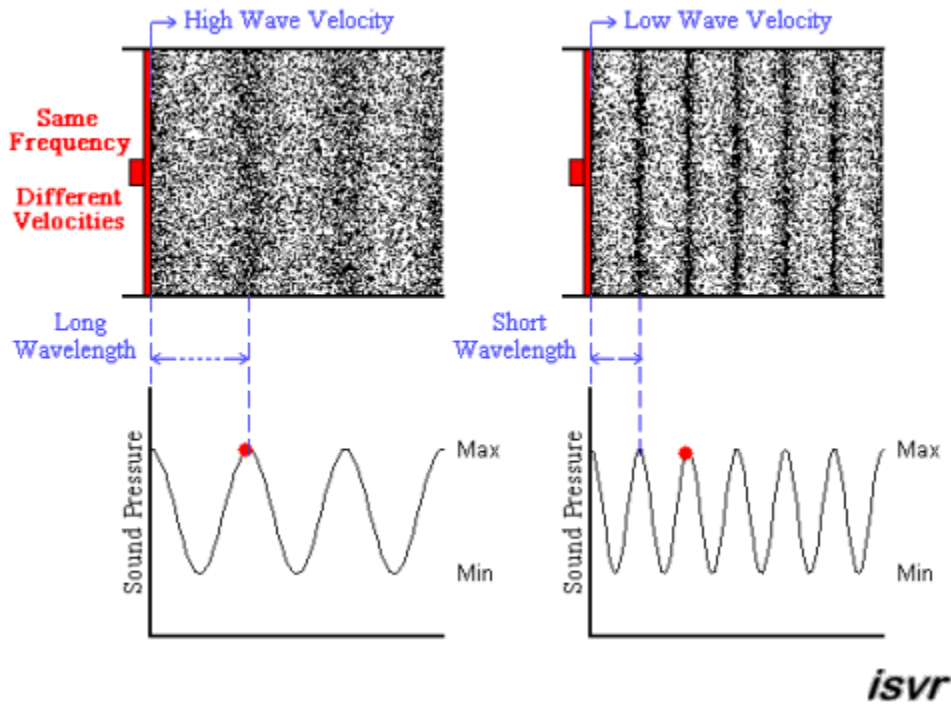
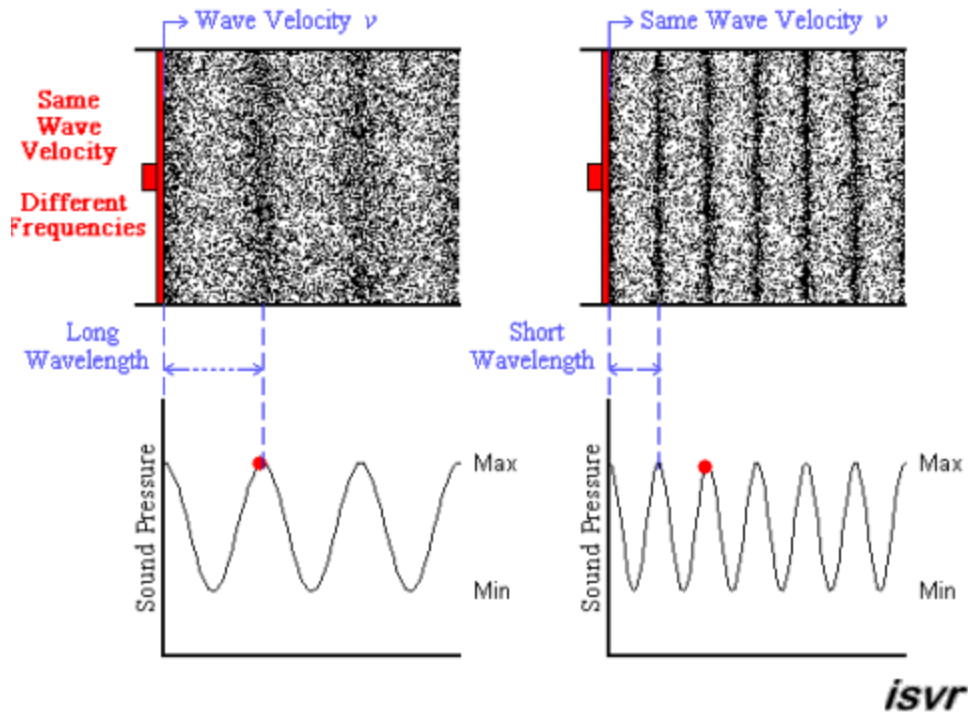
Light wave



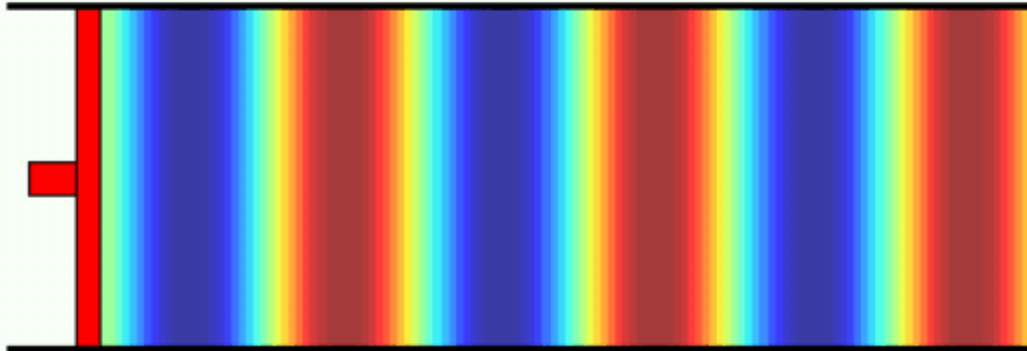
isvr

Acoustic Longitudinal Wave

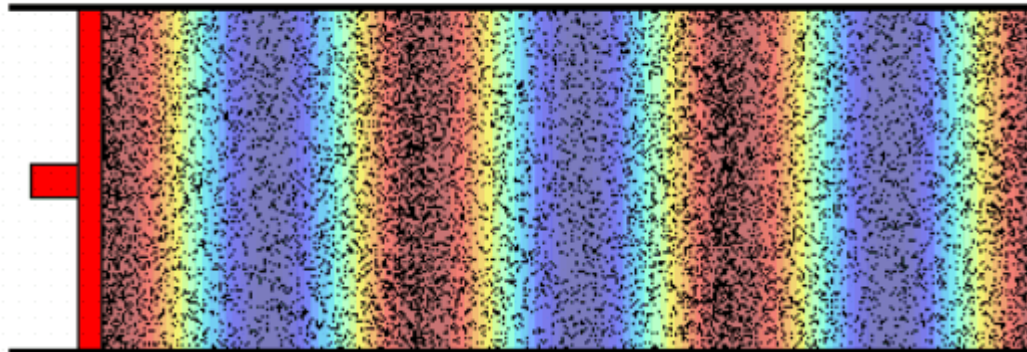




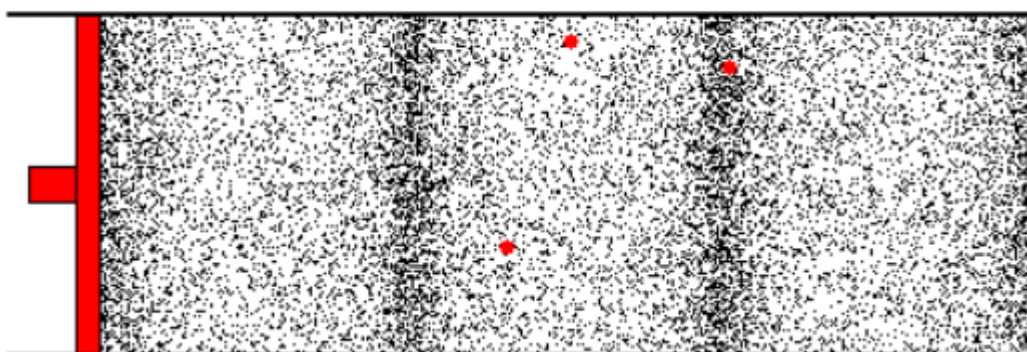
Longitudinal Wave

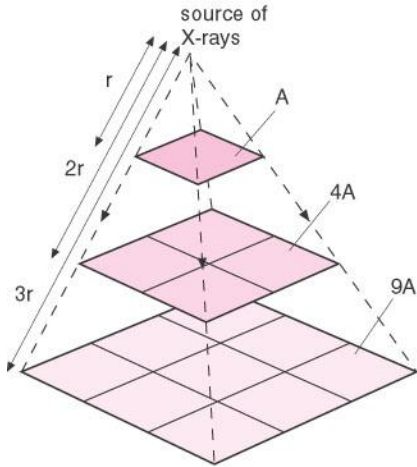


Longitudinal Wave



Longitudinal Wave





Mode (0,1)



Mode (1,1)



Mode (2,1)



Mode (0,2)



Mode (1,2)

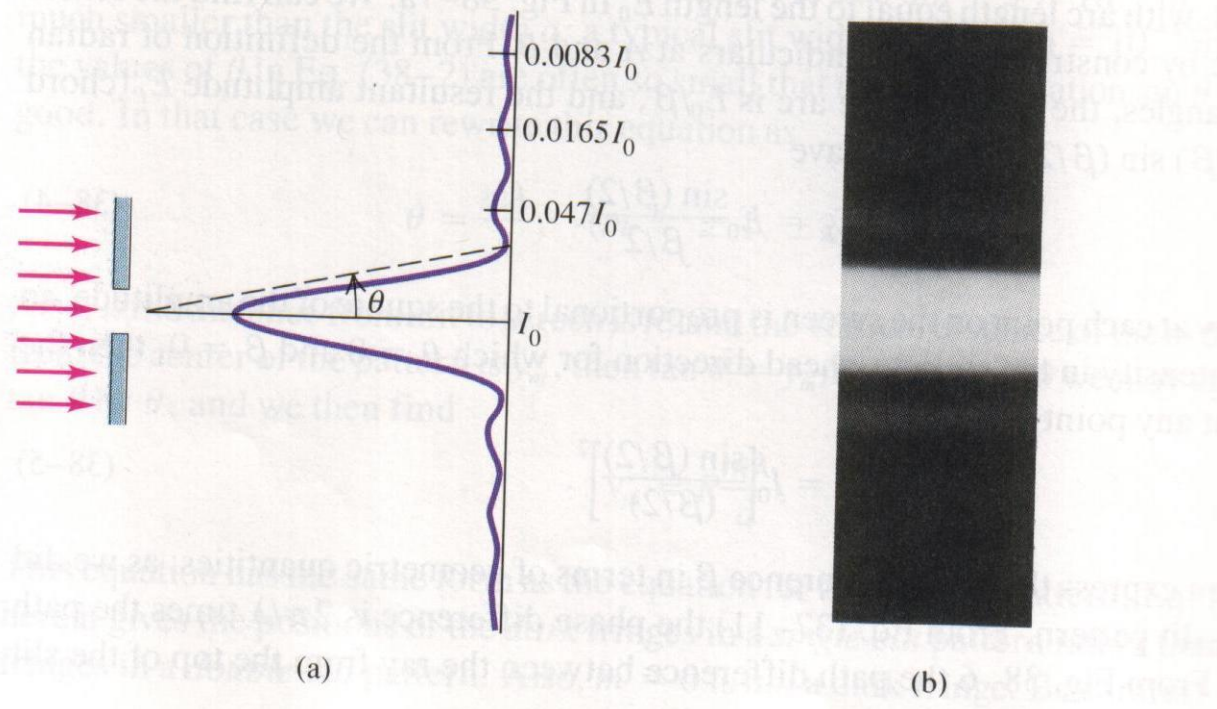
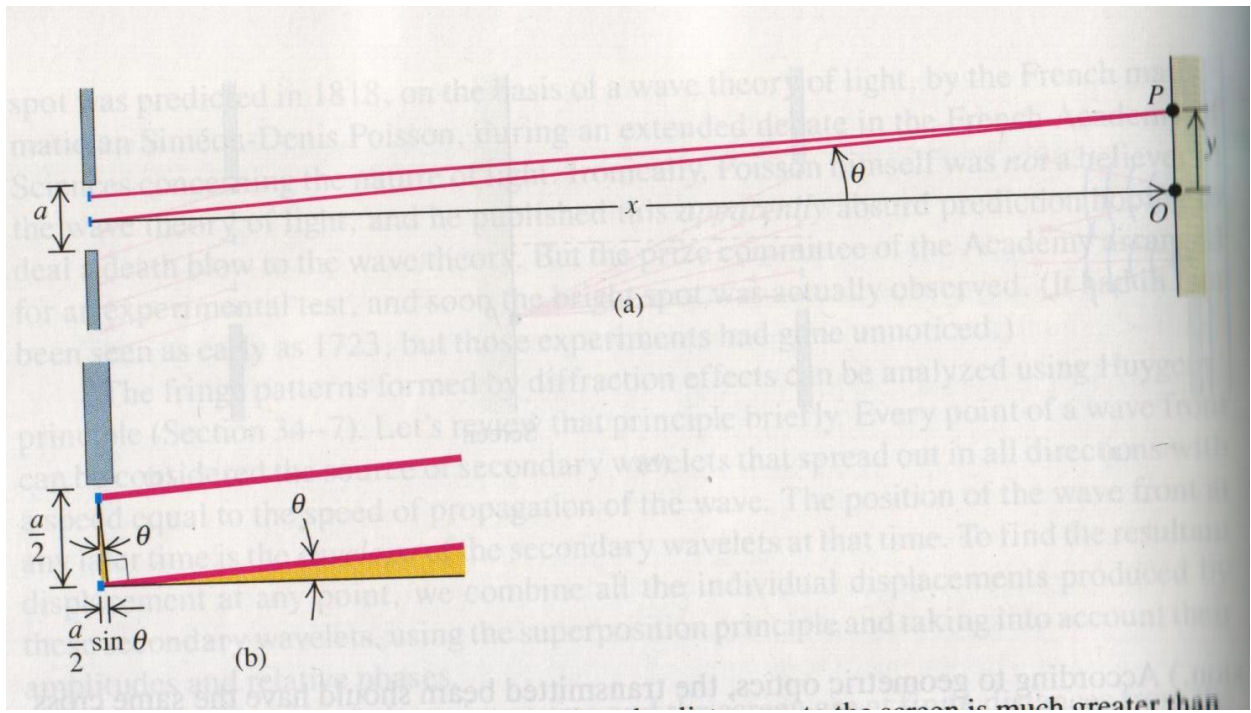


Mode (0,3)

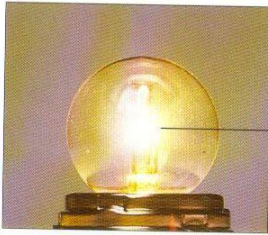


isvr





SOURCES OF LIGHT

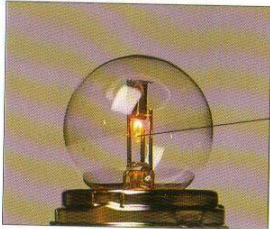


BRIGHT FILAMENT LAMP

This spectrum shows which colours are produced



All colours of light together combine to produce white
BRIGHT FILAMENT LAMP
 With a high **electric current**, the whole spectrum of visible light is produced (see p. 59).



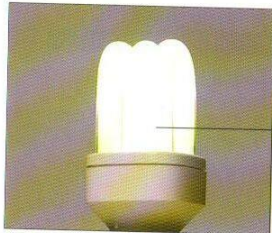
DIM FILAMENT LAMP

Red, yellow, and green light combine to produce orange



Lamp appears orange *No blue light produced*

DIM FILAMENT LAMP
 With a smaller current, the **temperature** of the **filament** (see pp. 52-53) is low.



FLUORESCENT LAMP

Lamp produces certain colours in each part of the spectrum



All three types of cone are stimulated and lamp appears white

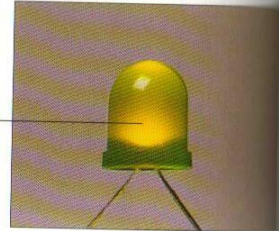
FLUORESCENT LAMP
 In a **fluorescent lamp**, chemicals called **phosphors** produce colours in many parts of the spectrum.

LED produces colours in the green part of the spectrum



LED appears green

GREEN LED
 An LED (light-emitting diode) is made of a **semiconductor**, and produces certain colours of light.



GREEN LED

Two colours of light very close together in the orange part of the spectrum are produced



Lamp appears orange

SODIUM LAMP
 In a sodium lamp, an electric current excites electrons in sodium vapour, giving them extra energy. The electrons give the energy out as light.



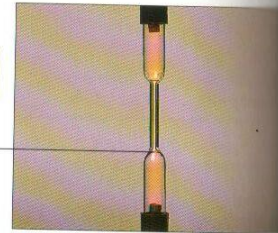
SODIUM LAMP

Only certain colours characteristic of neon are produced



Lamp appears orange

NEON TUBE
 In a similar way to a sodium lamp, a neon discharge lamp produces a characteristic orange glow.



NEON TUBE

STANDARD CANDLES

Cepheid variables

Pulse \propto size

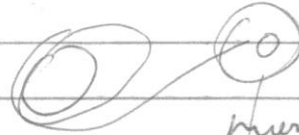
Pulse rate

↳ size of star

Look at size it appears on earth

↓
determine distance

Type I supernova



Draws matter from nearby source.

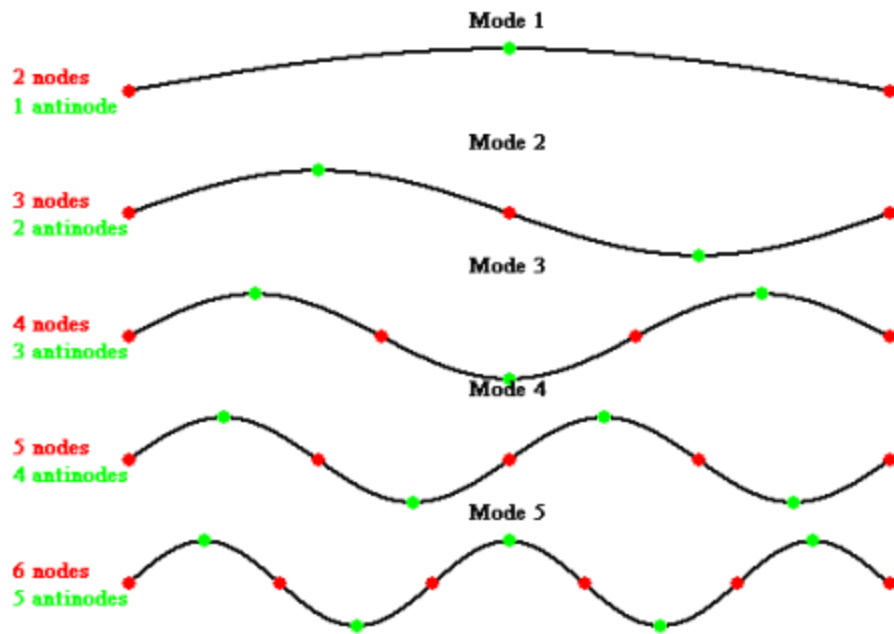
mass $\approx 1\frac{1}{2}$ sun

explodes

constant flash

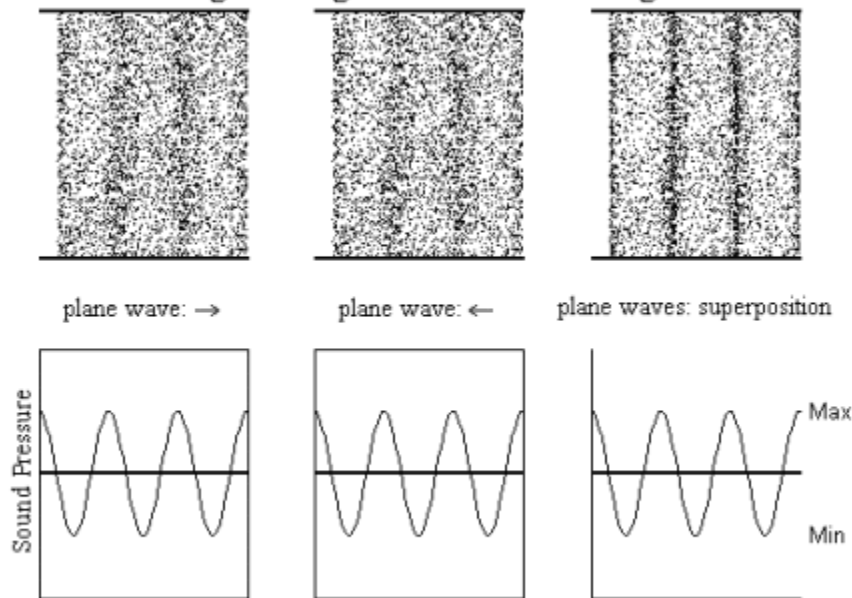
↓
determine distance

Closer objects → Parallax



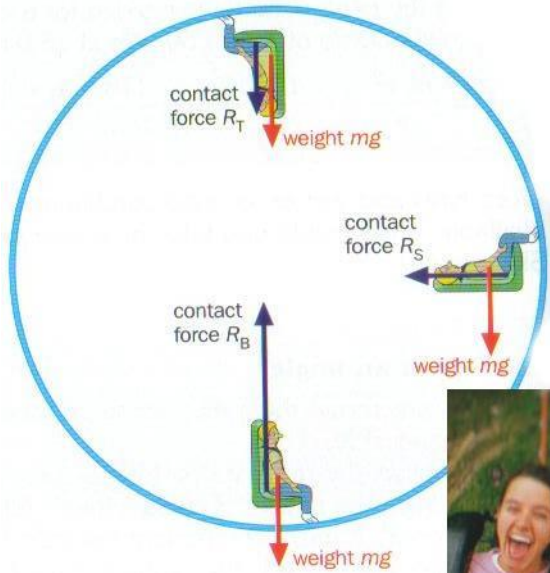
isvr

Creating Standing Waves from Travelling Waves



isvr

$$R_T + mg = \frac{m v^2}{r}$$



$$R_S = \frac{m v^2}{r}$$

$$R_B - mg = \frac{m v^2}{r}$$

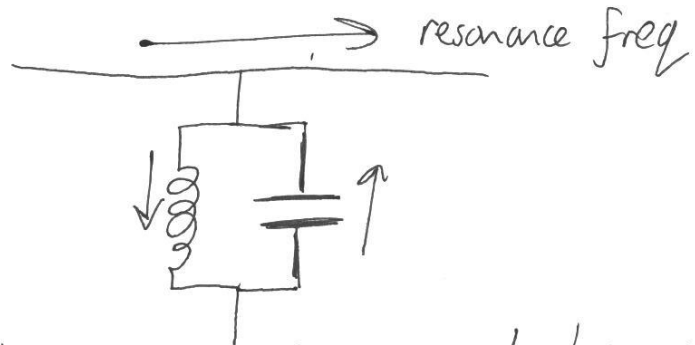


When

$$X_C = X_L$$

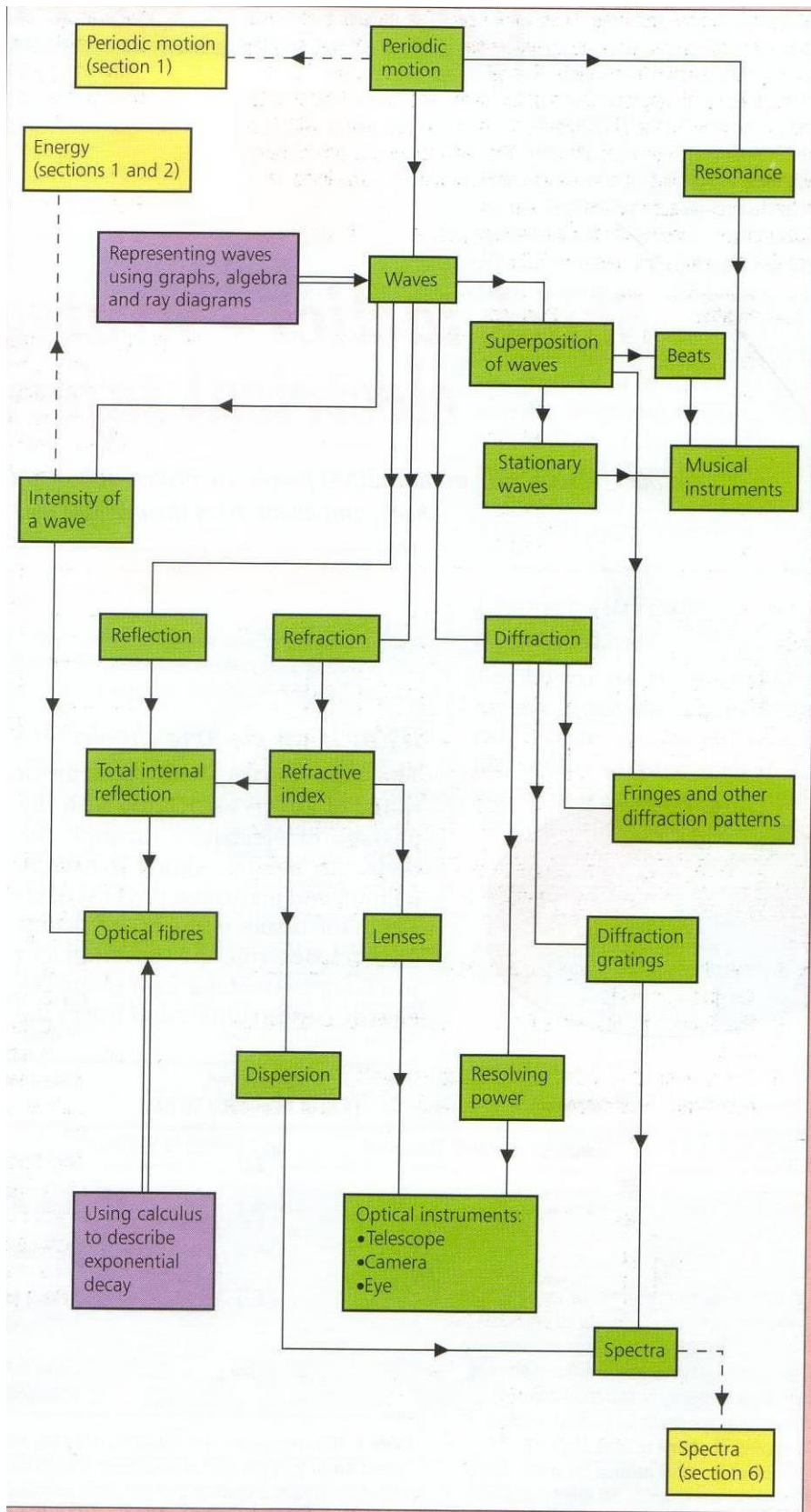
The AC resistance is the same

so current does not
Go to earth



$$\frac{1}{2\pi f_0 C} = 2\pi f_0 L$$

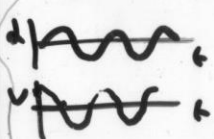
$$f_0^2 = \frac{1}{4\pi^2 LC} \Rightarrow f_0 = \frac{1}{2\pi\sqrt{LC}}$$



Spring
pendulum

KE \rightarrow PE
KE + PE = const

SHM



acceleration \propto -displacement

$a \propto -x$

$a = -\omega^2 x$

$F = -kx$

$x = x_0 \cos \omega t$

oscillations
↑
waves

centripetal acceleration $\rightarrow \frac{v^2}{r}$

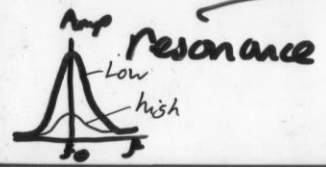
angular speed $\omega = \frac{\theta}{t} = \frac{2\pi}{T}$

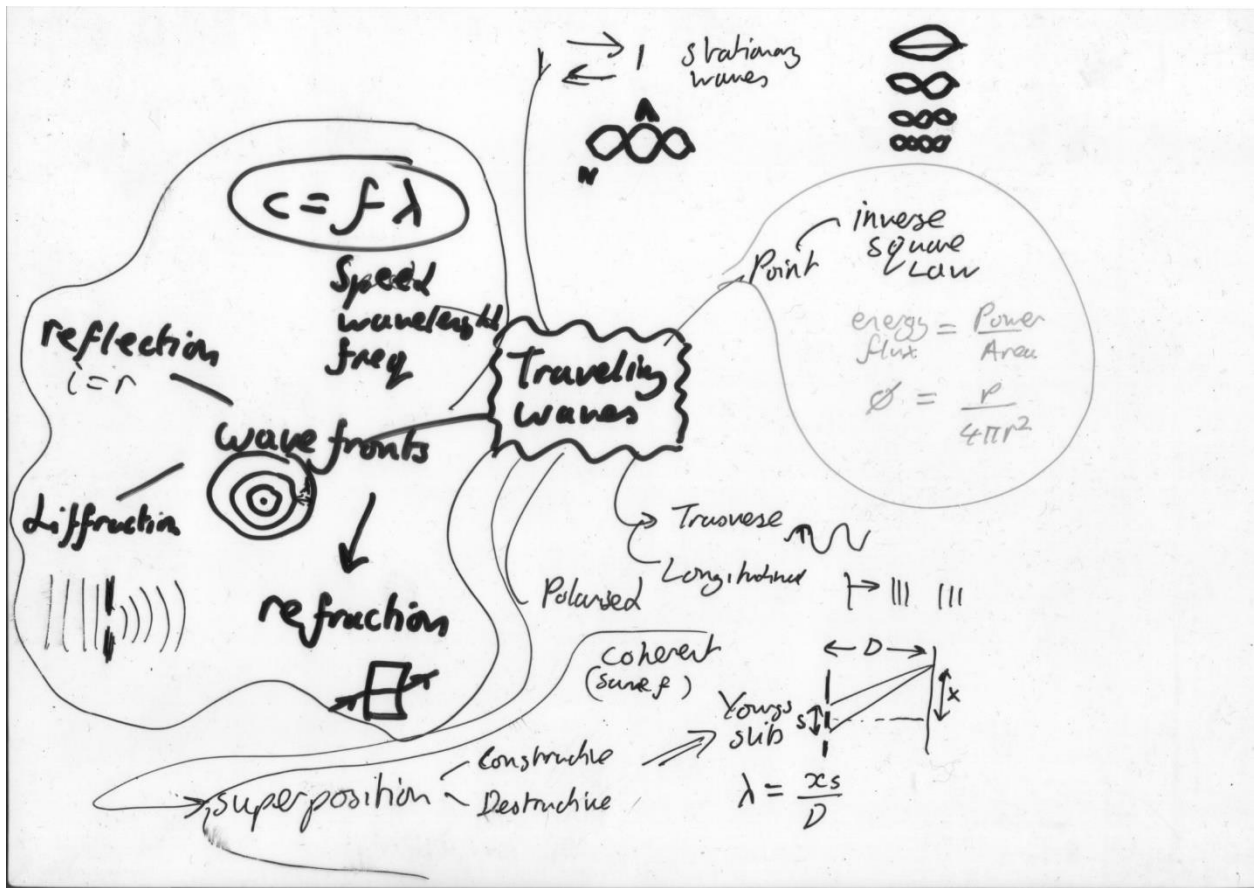
$v = r\omega$

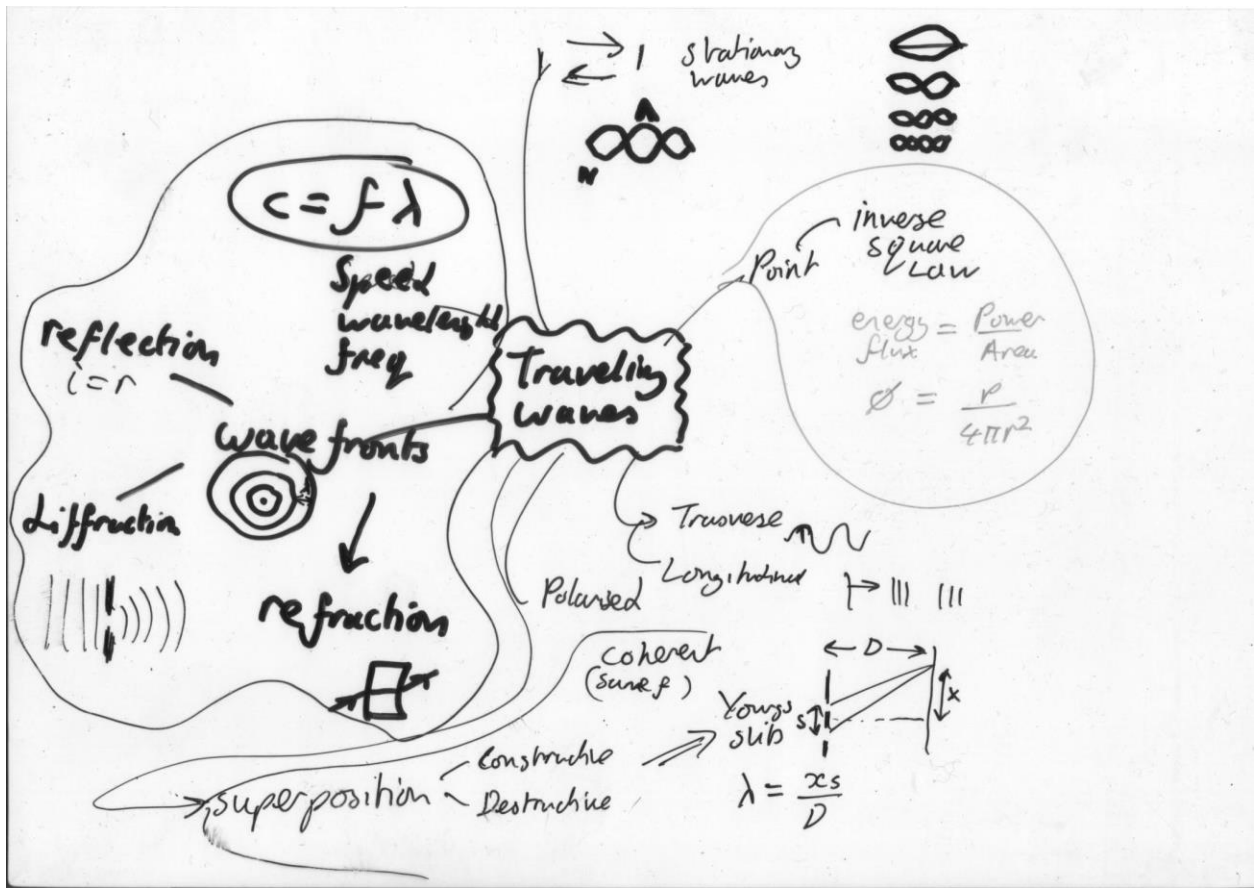
Conical
Pendulum

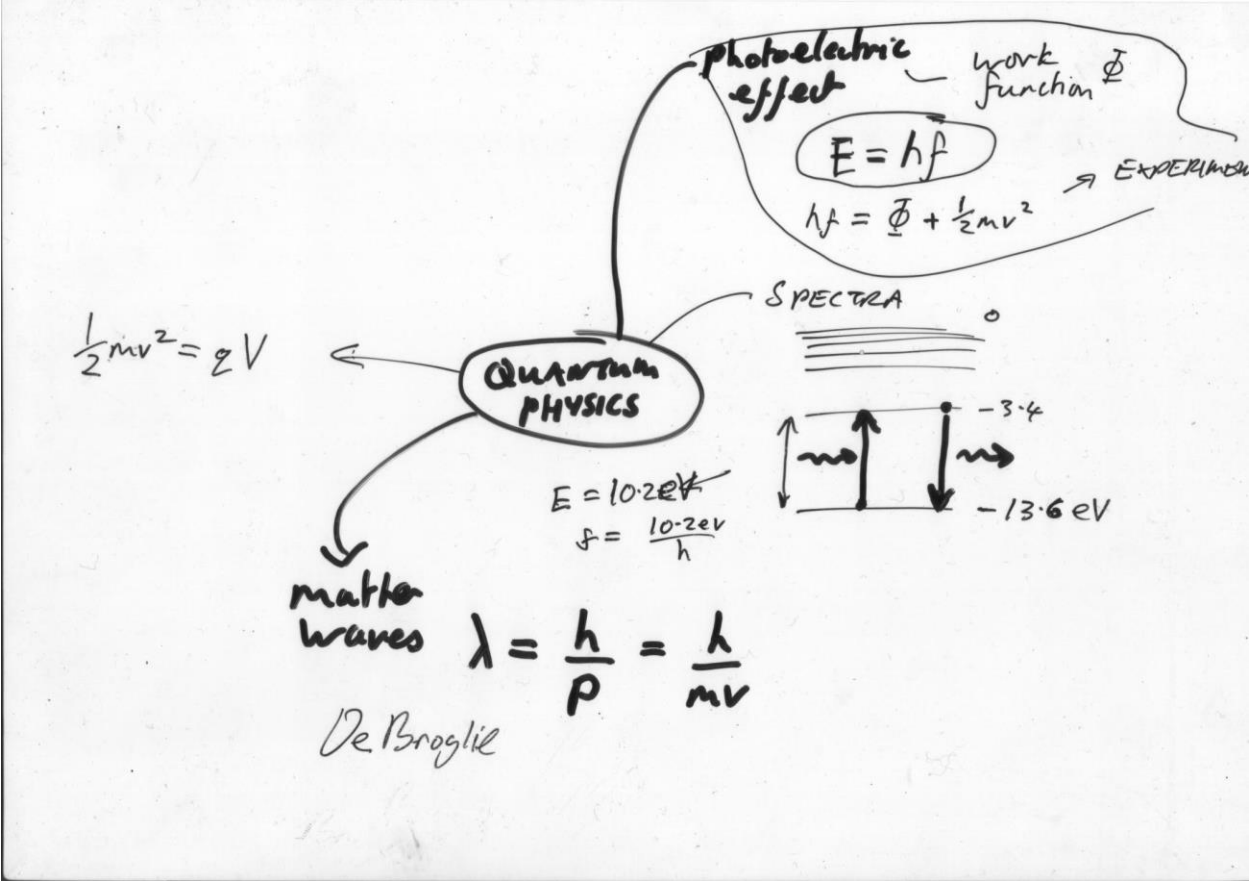


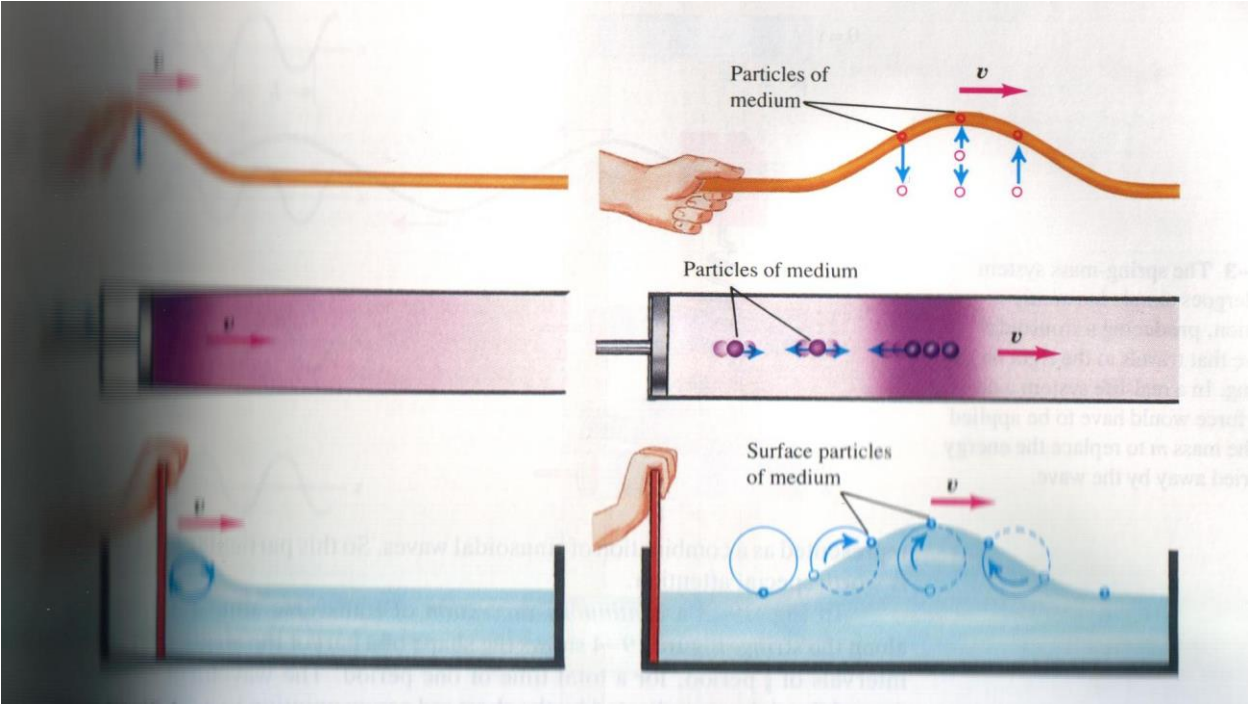
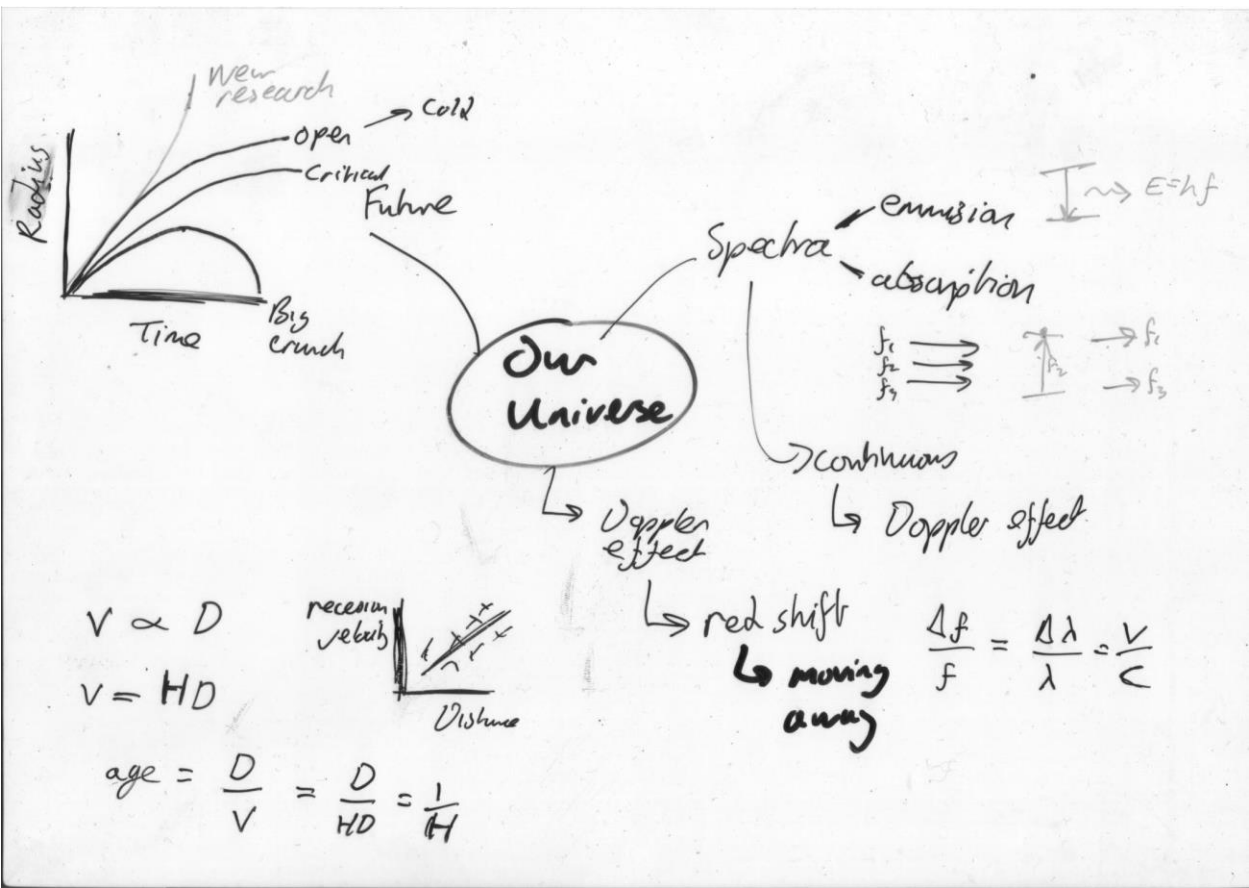
Periodic
motion $T = \frac{1}{f}$

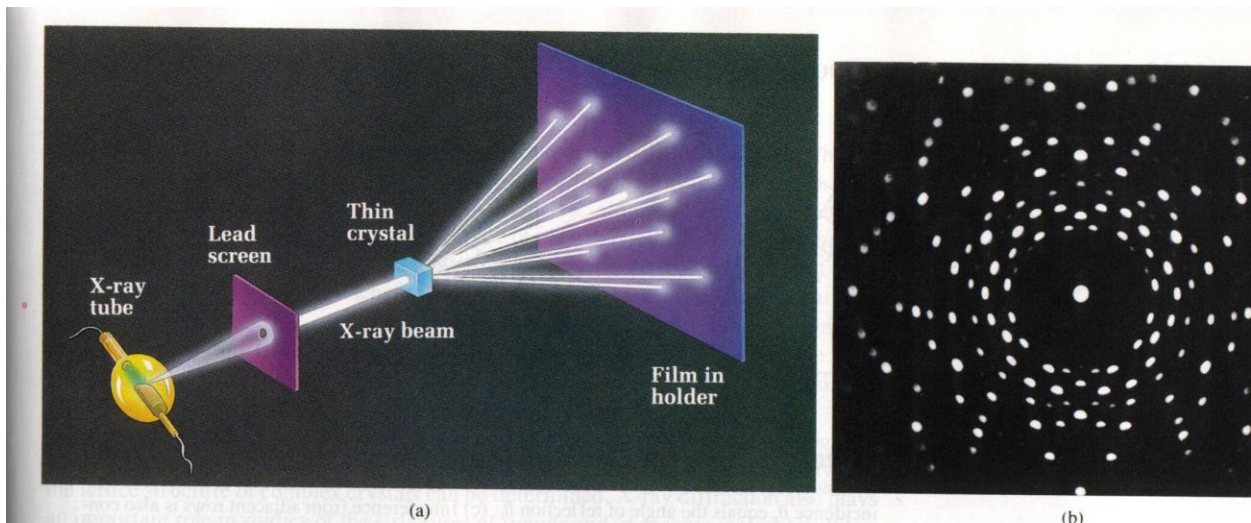




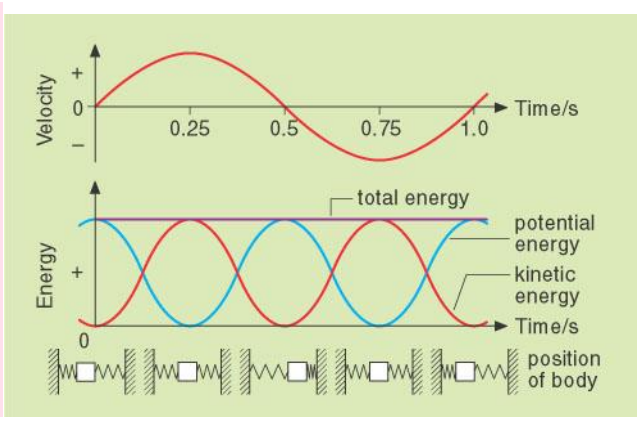
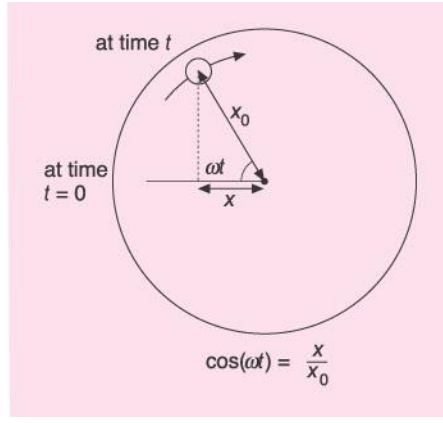
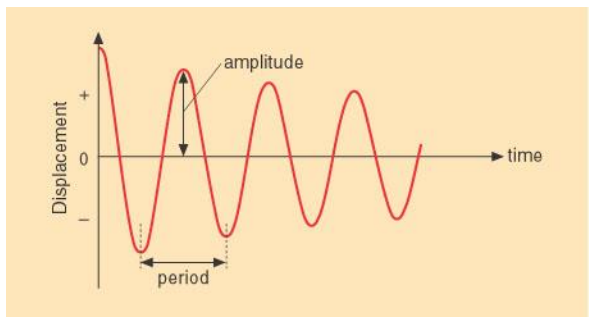
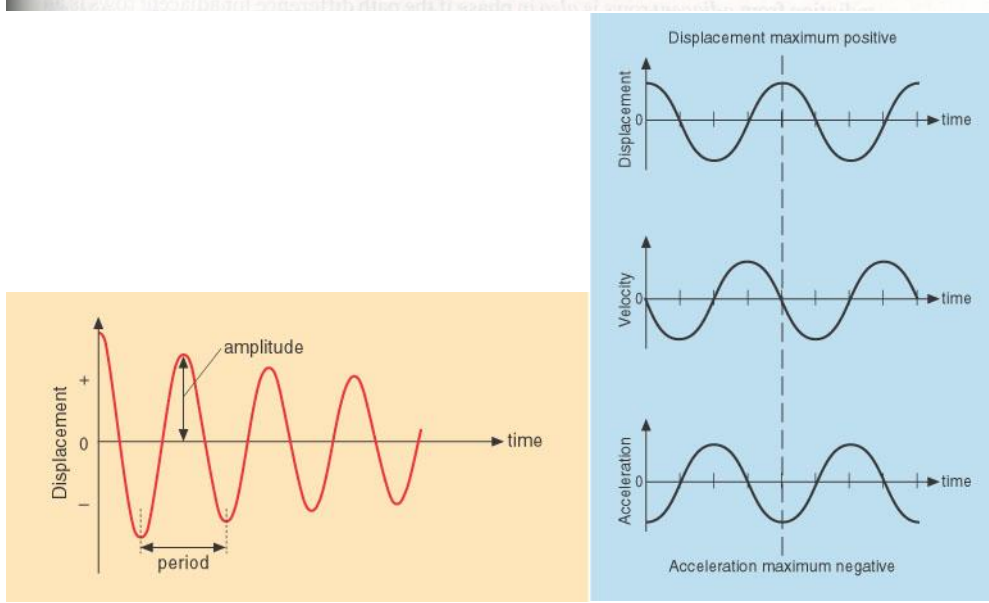


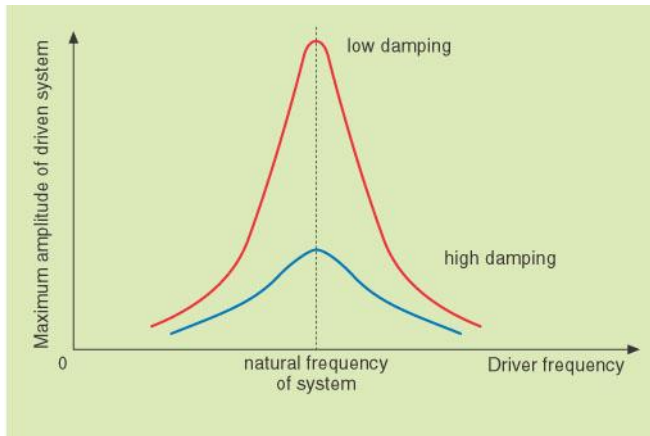






38-18 (a) In an x-ray diffraction experiment, most x rays pass straight through the crystal, but some are scattered, forming an interference pattern that exposes the film in a pattern related to the atomic arrangement in the crystal. (b) Laue diffraction pattern formed by directing a beam of x rays at a thin section of quartz crystal.



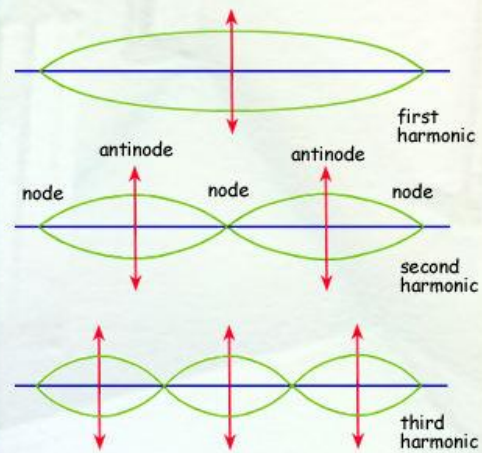


Standing or stationary waves are produced when two waves with the same frequency and amplitude move towards each other and become superimposed. The result is the creation of fixed nodal points of zero displacement which alternate with fixed antinodal points of maximum displacement. These nodes and antinodes are half a wavelength apart. Standing waves do not transfer energy along their length, as do progressive waves.

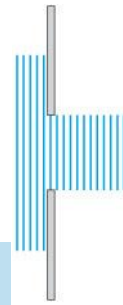
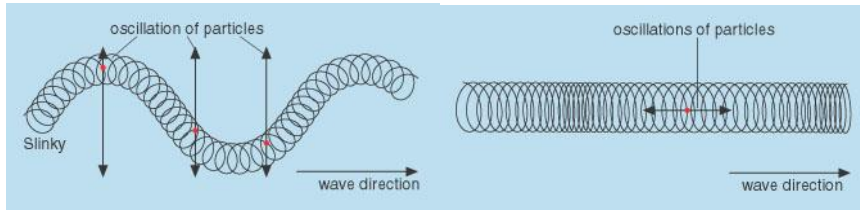
Standing Waves

Plucking a stretched string fixed at both ends, such as one on a stringed instrument, produces a standing wave with two fixed nodes at each end. The three simplest standing waves or modes are the first, second and third harmonics.

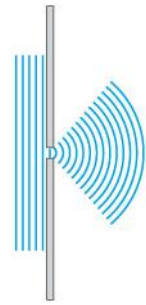
Standing waves such as the vibrations (overtones) of a stringed instrument can generate progressive waves. The progressive sound waves passing through air have the same frequency as the plucked string which generates them. The frequencies of each harmonic are whole number multiples of the lowest fundamental frequency. For a stringed instrument, when plucked, the string vibrates with all these frequencies.



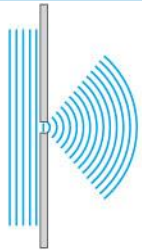
The first harmonic (or fundamental) has a node separation of $\frac{\lambda}{2} = L$
 The string length is half a wavelength.
 The second harmonic (or first overtone) has a node separation of $\frac{\lambda}{2} = \frac{L}{2}$
 The string length is one wavelength.
 The third harmonic (or second overtone) has a node separation of $\frac{\lambda}{2} = \frac{L}{3}$
 The string length is one and a half wavelengths.



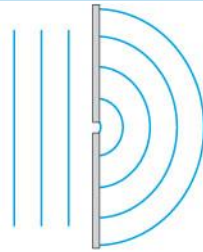
small wavelength, large gap



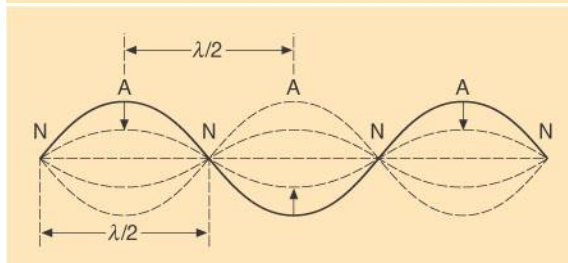
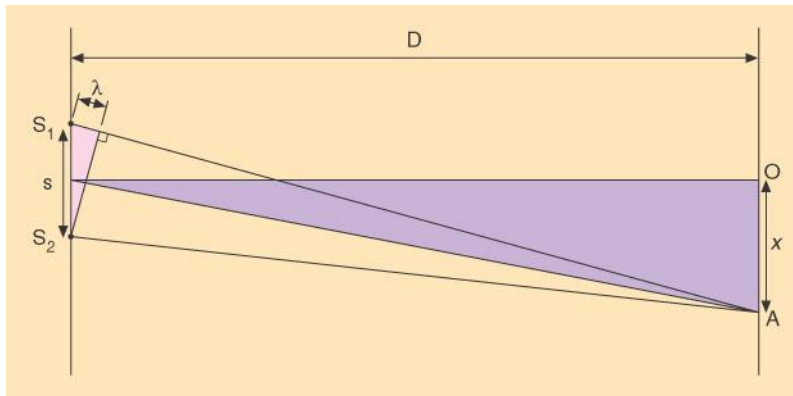
small wavelength, small gap

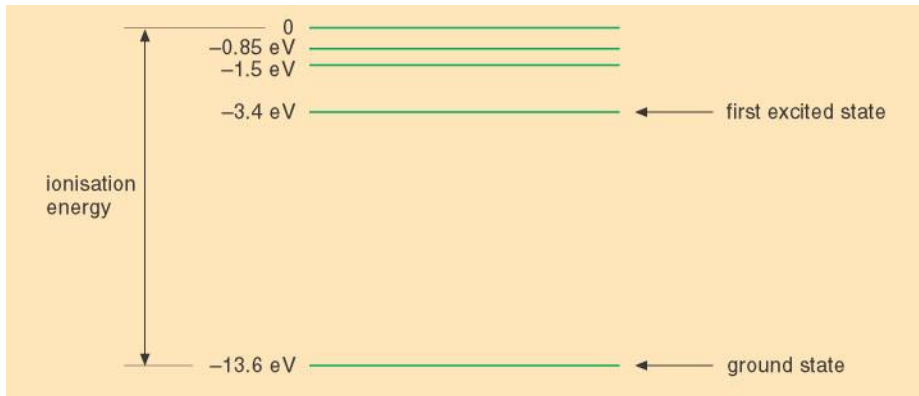


small wavelength, small gap



large wavelength, small gap

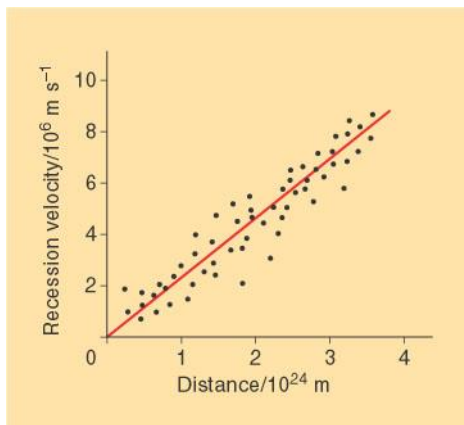




(a) Absorption bands in light from nearby galaxies



(b) Absorption bands in light from distant galaxies

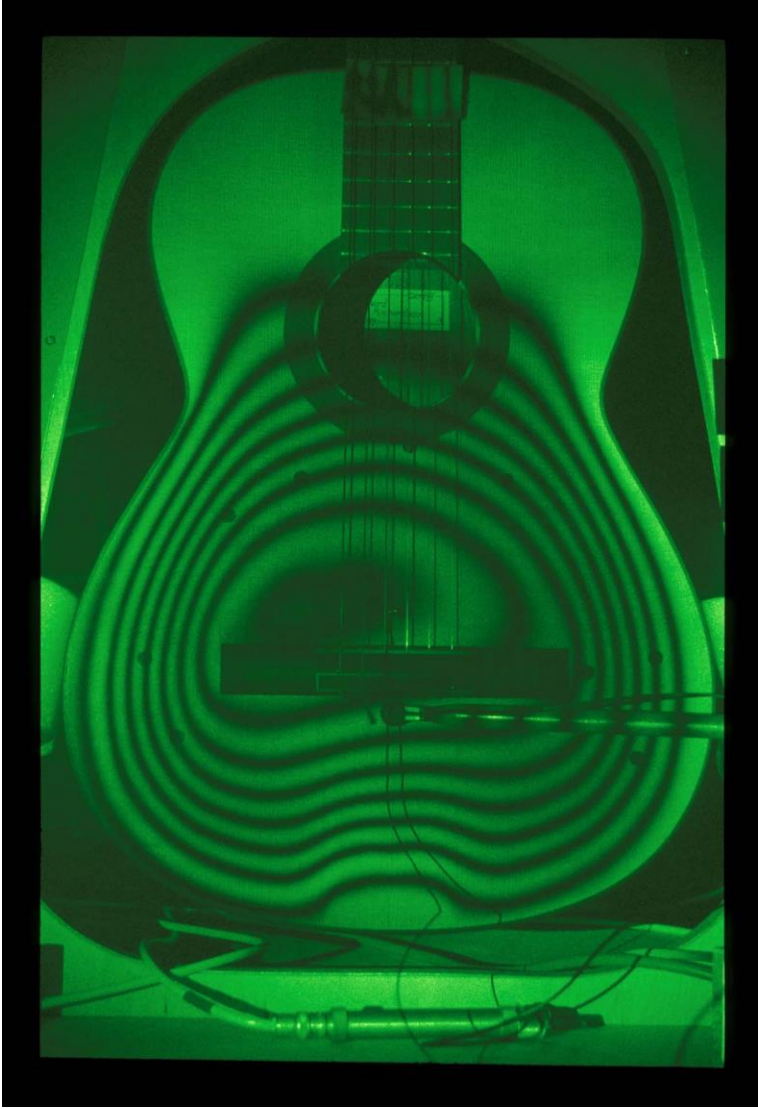


(a) Absorption bands in light from nearby galaxies



(b) Absorption bands in light from distant galaxies





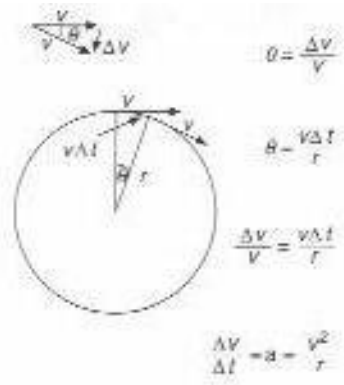
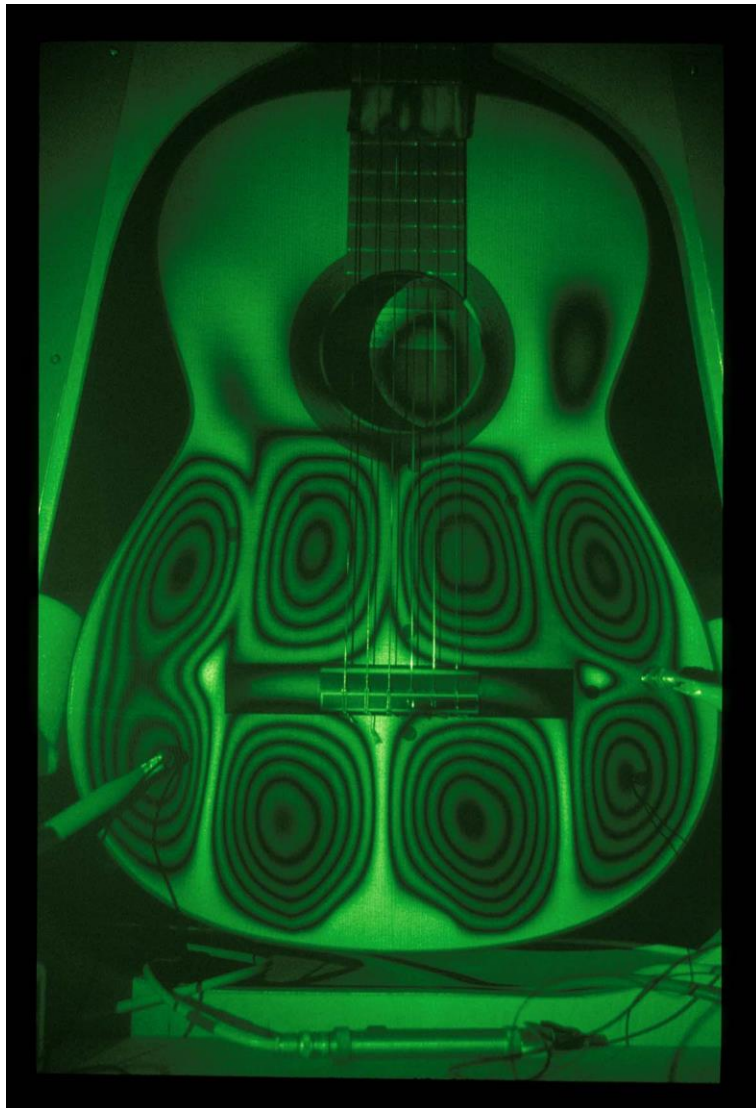
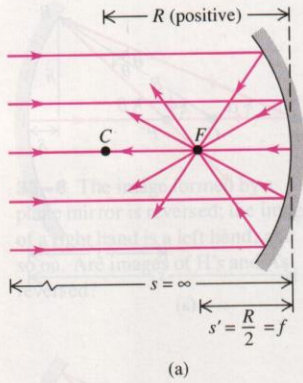
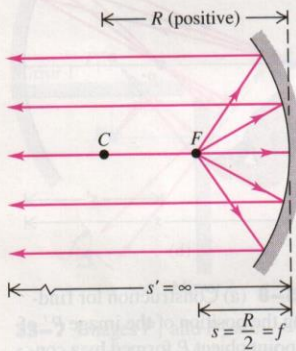


Figure 2.1 Centripetal acceleration = $\frac{v^2}{r}$



(a)



(b)

35-9 (a) Incident rays parallel to the axis converge to the focal point F of a concave mirror. (b) Rays diverging from the focal point F of a concave mirror are parallel to the axis after reflection. The angles are exaggerated for clarity.

Focal Point

When the object point P is very far from the mirror ($s = \infty$), the incoming rays are parallel. From Eq. (37-4) the image distance s' is given by

$$\frac{1}{\infty} + \frac{1}{s'} = \frac{2}{R}, \quad s' = \frac{R}{2}.$$

The situation is shown in Fig. 35-9a. A beam of incident parallel rays converges, after reflection, to a point F at a distance $R/2$ from the vertex of the mirror. Point F is called the **focal point**, and its distance from the vertex, denoted by f , is called the **focal length**. We see that f is related to the radius of curvature R by

$$f = \frac{R}{2}.$$

We can discuss the opposite situation, shown in Fig. 35-9b. When the image distance s' is very large, the outgoing rays are parallel to the optic axis. The object distance s is then given by

$$\frac{1}{s} + \frac{1}{\infty} = \frac{2}{R}, \quad s = \frac{R}{2}.$$

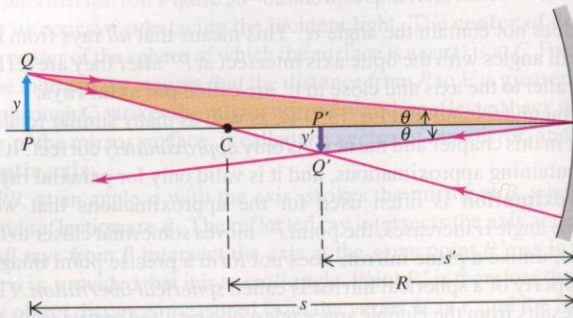
In Fig. 35-9b, rays coming to the mirror from the focal point are reflected parallel to the optic axis. Again we see that $f = R/2$.

Thus the focal point F of a concave spherical mirror has the properties that (1) an incoming ray parallel to the optic axis is reflected through the focal point and (2) an incoming ray that passes through the focal point is reflected parallel to the optic axis. In spherical mirrors these statements are true only for paraxial rays; for parabolic mirrors they are *exactly* true.

We will usually express the relationship between object and image distances for a mirror, Eq. (35-4), in terms of the focal length f :

$$\frac{1}{s} + \frac{1}{s'} = \frac{1}{f}.$$

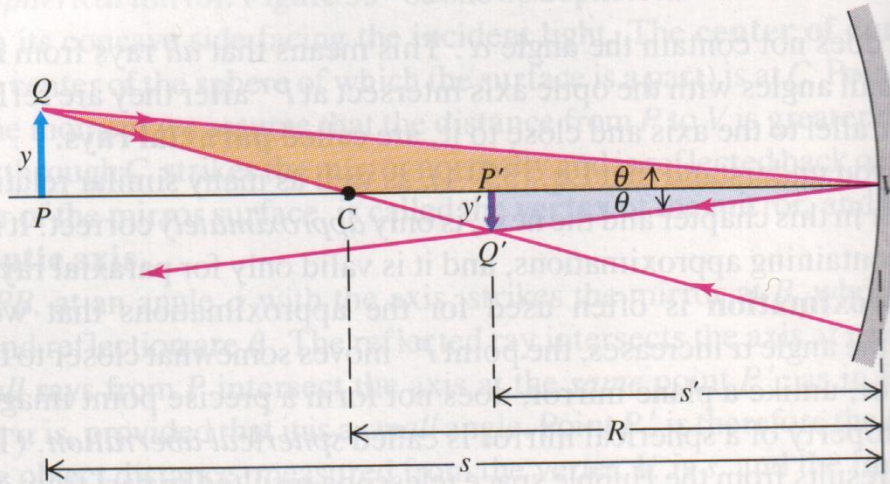
Now suppose we have an object with finite size, represented by the arrow PQ in Fig. 35-10, perpendicular to the axis PV . The image of P formed by paraxial rays is P' . The object distance for point Q is very nearly equal to that for point P , so the image of Q



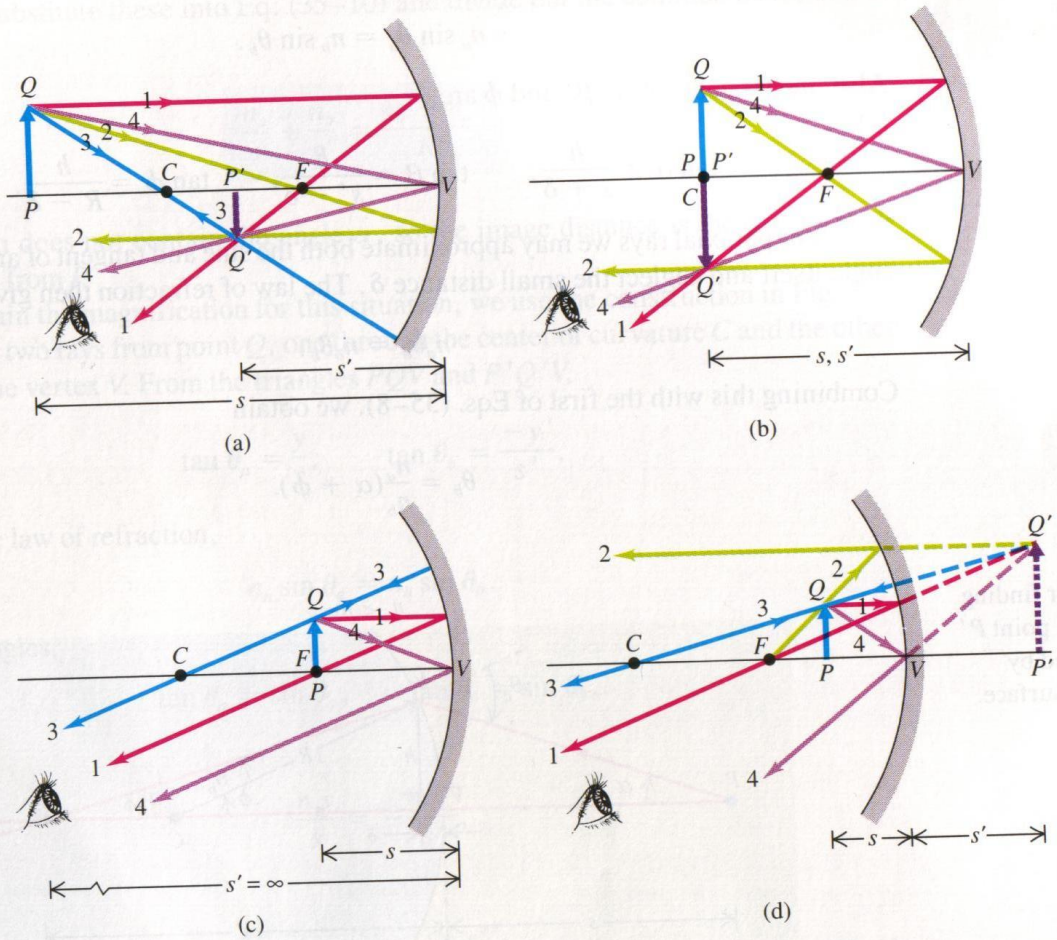
35-10 Construction for determining the position, orientation, and height of an image formed by a concave spherical mirror.

$$\frac{1}{s} + \frac{1}{s'} = \frac{1}{f}$$

Now suppose we have an object with finite size, represented by the arrow PQ in Fig. 35-10, perpendicular to the axis PV . The image of P formed by paraxial rays is P' . The object distance for point Q is very nearly equal to that for point P , so the image



35-10 Construction for determining the position, orientation, and height of an image by a concave spherical mirror.



35-17 Image of an object at various distances from a concave mirror, showing principal rays.

Particles as waves

de Broglie's hypothesis

a particle with momentum p has an associated wavelength λ given by

$$\lambda = h/p$$

this wavelength is called the *de Broglie wavelength*

- Because h is very small, only atomic-sized particles will have sufficiently small momentum for λ to be detectable.
- If the de Broglie wavelength for an electron is comparable to atomic dimensions then *electron diffraction* effects can be observed (principle of electron microscope).
- If the de Broglie wavelength for an electron is comparable to *nuclear* dimensions (very high energy) then diffraction by the nucleus can be observed giving information on nuclear size.

Example 1

What is the de Broglie wavelength for an electron accelerated through a p.d. of 5.0 kV?

$$m_e = 9.1 \times 10^{-31} \text{ kg}$$

Stage 1

Apply the Law of Conservation of Energy to the accelerated motion:

Method 1

$$\frac{1}{2} mv^2 = eV$$

$$v = \sqrt{(2eV/m)}$$

$$= \sqrt{(2 \times 1.6 \times 10^{-19} \text{ C} \times 5000 \text{ V} / 9.1 \times 10^{-31} \text{ kg})}$$

$$= 4.2 \times 10^7 \text{ m s}^{-1}$$

$$p = mv = 3.8 \times 10^{-23} \text{ kg m s}^{-1} \text{ (or N s)}$$

Method 2

$$E_k = p^2/2m \text{ and } E_k = eV$$

$$p = \sqrt{(2meV)}$$

$$= 3.8 \times 10^{-23} \text{ kg m s}^{-1} \text{ (or N s)}$$

Stage 2

Apply the de Broglie relation:

$$\lambda = h/p$$

$$= 6.6 \times 10^{-34} \text{ J s} / 3.8 \times 10^{-23} \text{ N s}$$

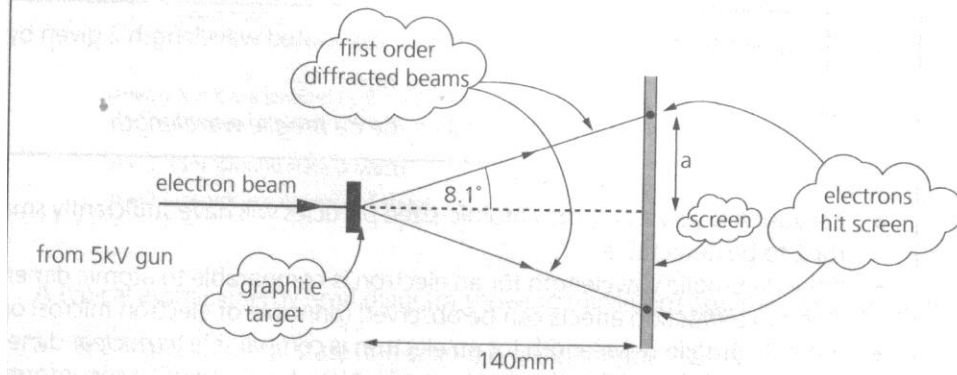
$$= 1.7 \times 10^{-11} \text{ m}$$

Example 2

The electron beam of example 1 is fired in a vacuum through a thin slice of graphite. The regularly spaced rows of carbon atoms make the crystal lattice behave like a diffraction grating to the electron beam, the diffraction following the ordinary rule for waves ($n\lambda = s \sin \theta_n$) using the de Broglie wavelength. If the atomic spacing is 1.2×10^{-10} m at what angle does the 1st order beam emerge?

The de Broglie wavelength has already been found to be
 1.7×10^{-11} m
so using the grating relation
 $1 \times (1.7 \times 10^{-11} \text{ m}) = 1.2 \times 10^{-10} \text{ m} \times \sin \theta_1$
 $\theta_1 = 8.1^\circ$

If it hits a screen 140 mm from the graphite describe what is seen on the screen.



Electron diffraction by graphite

$$a = 140 \text{ mm} \times \tan 8.1^\circ = 20 \text{ mm (2 sig. fig.)}$$

Because the graphite sample will have many layers of atoms at all possible orientations (polycrystalline) there will be many diffracted beams at the angle 8.1° , forming an emerging cone of electrons. They hit the screen making a circle of radius 20 mm.

Wave-particle duality

This dual nature of matter and radiation – some phenomena best described on a wave basis using ideas of interference, diffraction and wavelength and others on a particle basis using momentum and kinetic energy – is called *wave-particle duality*. It forms the basis of Quantum Mechanics, the form of physical laws best suited to describe atomic and nuclear phenomena.

The wave nature of a moving particle can be applied to the electron in a hydrogen atom, where the electron is represented as a standing wave whose wavelength in the fundamental mode is twice the atom's diameter. This model, where the electron wave is thought of as representing the probability of finding an electron at particular distances, is remarkably successful at predicting the energy levels in hydrogen.